I. INTRODUCTION

In recent years ferroelectric materials have attracted a lot of attention for use in nonvolatile memory devices, in particular, ferroelectric random access memories (FeRAM).\textsuperscript{1–3} The leading candidate for this application is PbZr\textsubscript{1−x}Ti\textsubscript{x}O\textsubscript{3} (PZT) perovskite which has a high Curie temperature and large remanent polarization. However, PZT thin films exhibit serious polarization fatigue during electric field cycling. More recently, a significant breakthrough for controlling the fatigue problems was reported in which ferroelectric materials belonging to the layered perovskite family, such as SrBi\textsubscript{2}Ta\textsubscript{2}O\textsubscript{9} (SBT), show essentially no polarization fatigue with electric field cycling.\textsuperscript{4}

SrBi\textsubscript{2}Ta\textsubscript{2}O\textsubscript{9} is a bismuth-layered pseudo-perovskite oxide. The crystal structure consists of [Bi\textsubscript{2}O\textsubscript{3}]\textsuperscript{2+} layers and perovskite-type [SrTa\textsubscript{2}O\textsubscript{7}]\textsuperscript{2−} units with double TaO\textsubscript{6} octahedral layers. Neutron and electron diffraction studies revealed that orthorhombic distortion with noncentrosymmetric space group A\textsubscript{2}1 am is responsible for the displacive-type ferroelectric behavior.\textsuperscript{5,6} The atomic displacements along the a axis from corresponding positions in the paraelectric phase with a parent tetragonal (I\textsubscript{4}mm) structure occur at the Curie temperature, \( T_C = 608\) K\textsuperscript{7} and these displacements cause the spontaneous ferroelectric polarization in SBT.\textsuperscript{8}

The effects of particle size on the physical properties of ferroelectric materials, especially the nanocomposites of electronic devices, have been extensively investigated theoretically using transverse the Ising model\textsuperscript{7,8} and Landau phenomenological theory\textsuperscript{9,10} and experimentally by x-ray diffraction (XRD)\textsuperscript{11,12} and Raman scattering.\textsuperscript{13} Frey et al.\textsuperscript{14} studied the effect of particle size on the crystal structure of BaTiO\textsubscript{3}; Ishikawa, Yoshikawa, and Okada\textsuperscript{15} investigated the effect of particle size on the Curie temperature in PbTiO\textsubscript{3} (PT); and the study in size effects on dielectric properties of PbZr\textsubscript{1−x}Ti\textsubscript{x}O\textsubscript{3} (PZT) was reported by Huang et al.\textsuperscript{10} In this article we present an in situ Raman scattering study of size effect on the ferroelectric phase transition in SBT ultrathin particles.

II. EXPERIMENT

The mechanical activation process has been described in detail in our previous studies.\textsuperscript{11,15} The starting materials used in this work were commercially available SrO (99.9% in purity, Sigma-Aldrich Chemical Company, Inc.), Bi\textsubscript{2}O\textsubscript{3} (99.9% in purity, Acros Organics), and Ta\textsubscript{2}O\textsubscript{5} (99% in purity, Sigma-Aldrich Chemical Company, Inc.) After 30 h mechanical activation of the starting mixed oxides, the milled powders were subjected to a postannealing up to 950°C for 1 h with 5°C/min heating rate and 10°C/min cooling rate. The SBT bulk sample was prepared by sintering the 30 h milled powder at 1225°C for 2 h. Both the annealed powder and sintered bulk sample were characterized by x-ray diffraction (XRD, X’Pert, Philips). To study the temperature induced ferroelectric phase transition in SBT powder, in situ high temperature Raman experiments were carried out using a TMS 93 Linkam thermal stage capable of holding temperature over the range between 300 and 770 K. All micro-Raman spectra were measured in the backscattering geometry using a Jobin Yvon T64000 Raman spectrometer with an Olympus microscope attachment and equipped with a liquid-nitrogen-cooled charge coupled device (CCD) detector. The 514.5 nm line of an argon-ion laser was used as the excitation source.

III. RESULTS AND DISCUSSION

Figure 1 shows the XRD patterns of the samples annealed at various temperatures and the sintered bulk sample. The polycrystalline single phase SBT was formed in all annealed samples and this SBT phase was well established with
the increase of annealing temperatures.\textsuperscript{16} After fitting the XRD patterns with a Gaussian line shape, the average particle size was determined from the full width at half maximum (FWHM) of the strongest x-ray diffraction peak (115) using Scherrer equation which assumes the small crystallite size to be the only case of line broadening\textsuperscript{17}

\[ D_{\text{XRD}} = \frac{K\lambda}{B \cos \theta} \]

where \( D_{\text{XRD}} \) is the particle diameter, \( \lambda \) is the x-ray wavelength, \( B \) is the FWHM of the diffraction line, \( \theta \) is the angle of diffraction, and the constant \( K \approx 1 \). In order to remove the contribution from instrumental broadening, the FWHM of the (115) line in XRD trace of the sintered bulk sample was subtracted as described by Eq. (2)\textsuperscript{18}

\[ B^2 = a^2 - b^2 \]  

where \( B \) is the FWHM of the true diffraction profile, \( a \) and \( b \) are the measured FWHM of equivalent diffraction lines in the annealed powders and the sintered bulk sample, respectively. The calculated size of our annealed SBT powders ranges from 11 to 71 nm.

Figure 2 shows the temperature dependence of the low-frequency (10–300 cm\textsuperscript{-1}) Raman spectra of the bulk sample. The lowest frequency peak, \( A_1(z) \) mode shows an obvious redshift and broadening with increasing temperature, especially mostly softens near 560 K, and disappears above 580 K as shown in Fig. 2 inset. This typical soft mode behavior is responsible for the ferroelectric phase transition in bulk SBT sample.\textsuperscript{19} It is noted that the lowest frequency peak \( A_1 \) disappears below the Curie temperature (\( T_c = 589 \) K) obtained from the temperature dependence of dielectric constant (shown in Fig. 2 inset). A similar phenomenon has been reported by Kojima.\textsuperscript{19,20} Thus, the phase transition temperatures in SBT powders cannot be directly read from the \textit{in situ} Raman spectra as previous Raman scattering studies in Ba\textsubscript{3}TiO\textsubscript{5},\textsuperscript{21} Pb\textsubscript{3}TiO\textsubscript{5},\textsuperscript{13} and PbZr\textsubscript{1-x}Ti\textsubscript{x}O\textsubscript{3}.	extsuperscript{22}

Following Burns and scott\textsuperscript{23} and Kojima,\textsuperscript{19,20} we fitted the soft mode \( A_1 \) by a simple damped harmonic oscillator model after the correction of the Bose factor. According to the following two equations:

\[ \omega_p = \omega_{\text{TO}} \left[ 1 - \frac{1}{2} \left( \frac{\gamma}{\omega_{\text{TO}}} \right)^2 \right]^{1/2} \]  

\[ \frac{I(\omega_p)}{I(0)} = \frac{\omega_{\text{TO}}}{\gamma} \frac{1}{4} \left( \frac{\omega_{\text{TO}}}{\gamma} \right)^2 \]  

we measured the frequency corresponding to the peak intensity, \( \omega_p \), the intensity at the peak, \( I(\omega_p) \), and the intensity at zero frequency, \( I(0) \). A computer fitting with a Lorentzian line shape was carried out and the damping factors \( \gamma \) were obtained. Figure 3 shows the damping factors of the soft mode \( A_1(z) \) for the SBT bulk sample and nanoparticles as a function of temperature. The dark symbols represent the damping factors obtained from Eqs. (3) and (4). The solid lines are fitting results by the universal scaling theory which predicted the temperature dependence of damping factors \( \gamma(T) \) of the soft modes as\textsuperscript{24}

\[ \gamma(T) = \gamma(0) t^{-0.5} \]

where \( t = (T_c - T)/T_c \). The good agreements between the experimental data and the theoretical fitting results were ob-

\[ \text{FIG. 1. XRD patterns of sintered bulk sample and samples annealed at various temperatures.} \]

\[ \text{FIG. 2. Low-frequency Raman spectra of SBT bulk sample as a function of temperature. The inset shows the temperature dependencies of dielectric constant and frequency of soft mode of the SBT bulk sample.} \]

\[ \text{FIG. 3. Temperature dependence of damping factors of SBT bulk sample and nanoparticles of various sizes. The arrows indicate the transition temperature} \ T_c. \]
FIG. 4. Ferroelectric transition temperatures as a function of particle size.

observed in both the bulk sample and nanoparticles with various particle sizes. The estimated ferroelectric phase transition temperatures were also indicated by the arrows in Fig. 3.

Figure 4 shows the phase transition temperatures $T_c$ as a function of particle size. Similar to the previous study of size measurement, $25$ and $7$ nm by x-ray diffraction. $12$ It is even $20$ nm in this work. The solid line in Fig. 4 is the value size decreases, in particular when the particle size is less than $20$ nm in this work. The solid line in Fig. 4 is the value obtained by the following equation:$^{13}$

$$T_c(D) = T_c(\infty) - C/(D - D')$$

where $T_c(D)$ is the transition temperature of the particles with $D$ nm in size; $T_c(\infty)$ is the transition temperature of the bulk sample, which is taken to be $605$ K (see Fig. 3) although it is relatively higher than $T_c$ obtained from dielectric constant; $T_c$ of SBT nanoparticles was derived from Raman scattering, $C$ and $D'$ are fitting parameters. The above equation was first proposed by Ishikawa, Yoshikawa, and Okada$^{13}$ and the qualitative understanding based on the soft mode picture was given by Zhong et al.$^{25}$ According to previous definition,$^{13,25}$ a critical size $D_{\text{crit}}$ is the size at which $T_c = 0$ K. We obtained $D_{\text{crit}}$ of SBT from Eq. (6) to be $2.6$ nm. This value is much less than the values of PT, which is $13.8$ nm derived by Raman scattering,$^{13}$ $8.8$ nm by specific heat measurement,$^{25}$ and $7$ nm by x-ray diffraction.$^{15}$ It is even less than that of PZT ($4$ nm$^{10}$ and $4.2$ nm$^3$). This small critical size indicates another potential application of SBT as a promising ferroelectric material for the formation of low dimension and high density ferroelectric devices besides being fatigue free.

IV. CONCLUSION

The size effect on the ferroelectric phase transition in SrBi$_2$Ta$_2$O$_9$ nanoparticles has been studied by in situ Raman scattering. With the decrease of particle size, the transition temperature decreases dramatically. From an empirical expression, the critical size ($D_{\text{crit}} = 2.6$ nm) of SBT was obtained. Ferroelectricity will be lost when the particle size is less than $D_{\text{crit}}$. Compared with other ferroelectric materials, SBT has a relatively smaller critical size, which implies that SBT is a potential candidate for formation of ferroelectric devices of ultrafine size.